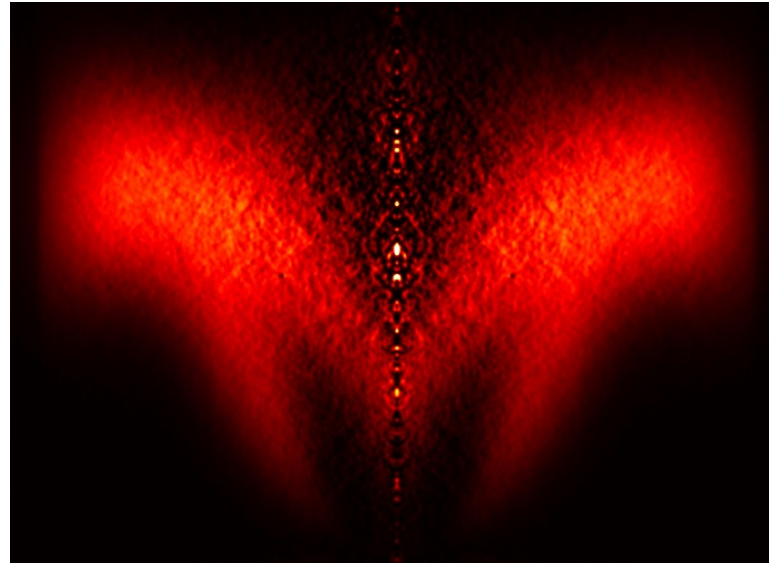


# *Calculation of flame describing functions from experimental data*

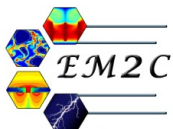


POLKA – 6<sup>th</sup> Scientific Workshop

17 March 2023

**Preethi Rajendram Soundararajan**

Department of Engineering, University of Cambridge, UK



# In this lecture

- Lecture:

  - Essentials of signal processing

- Exercise

  - Measuring FDFs experimentally

  - Calculating FDFs from experimental data

# Sampling theorem

How fast should you sample data?



# Sampling theorem

How fast should you sample data?

## Nyquist-Shannon criterion

An analog signal should be sampled at a rate **greater than twice** its highest frequency

$$f_s > 2 \times f_m$$

Sampling frequency  Maximum frequency 

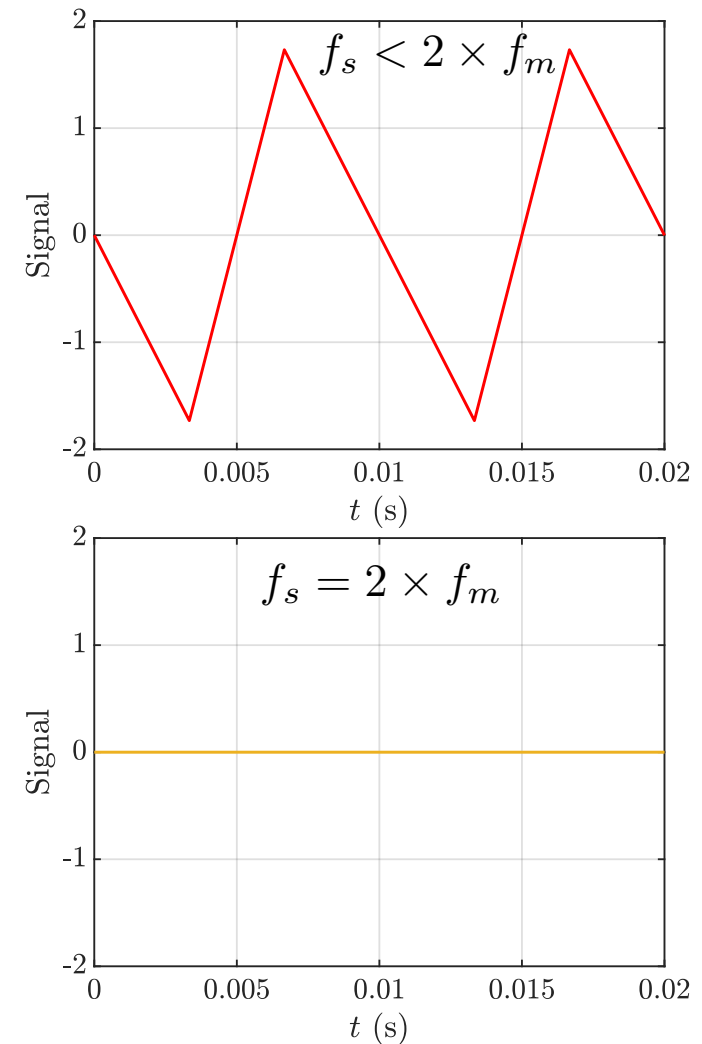
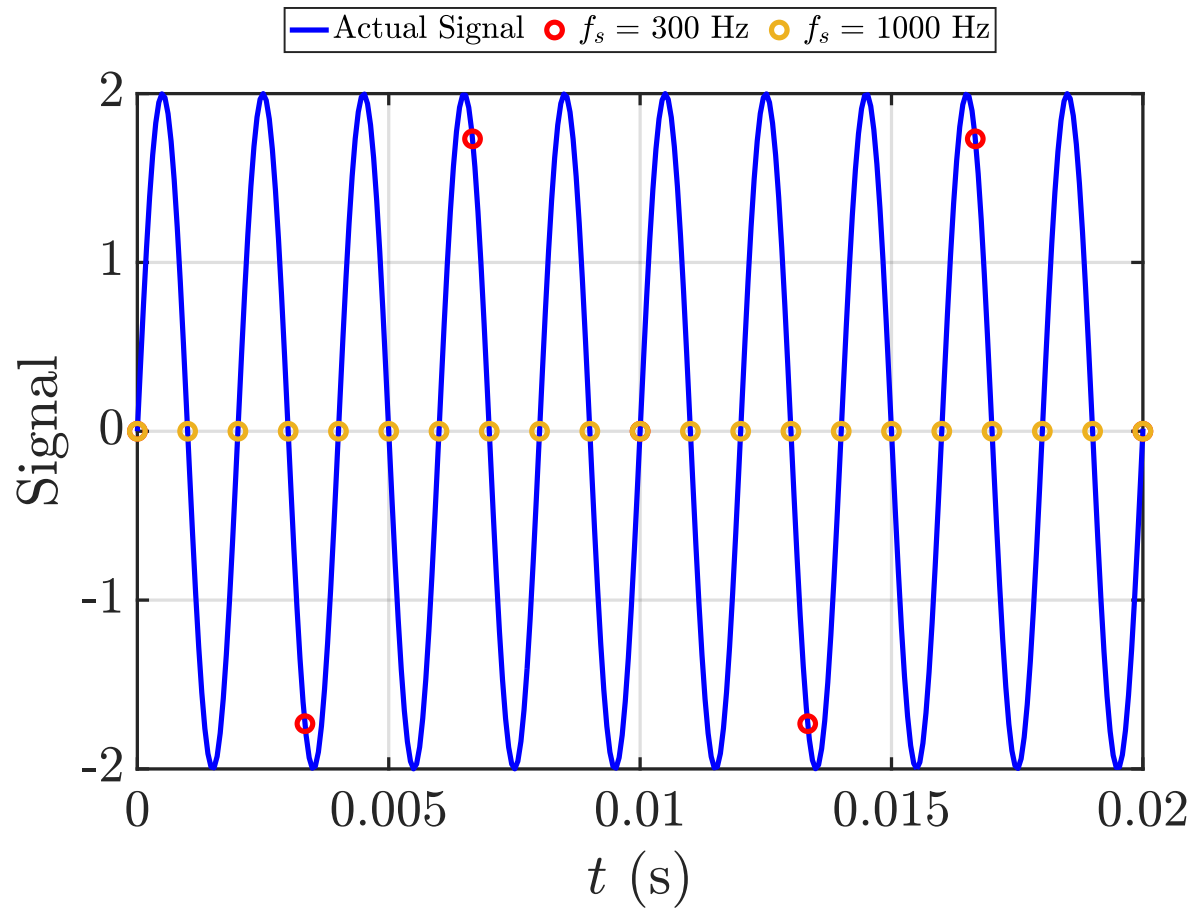
Nyquist frequency

$$f_{Nq} = 2 \times f_m$$

# Sampling theorem

Nyquist-Shannon criterion:  $f_s > 2 \times f_m$

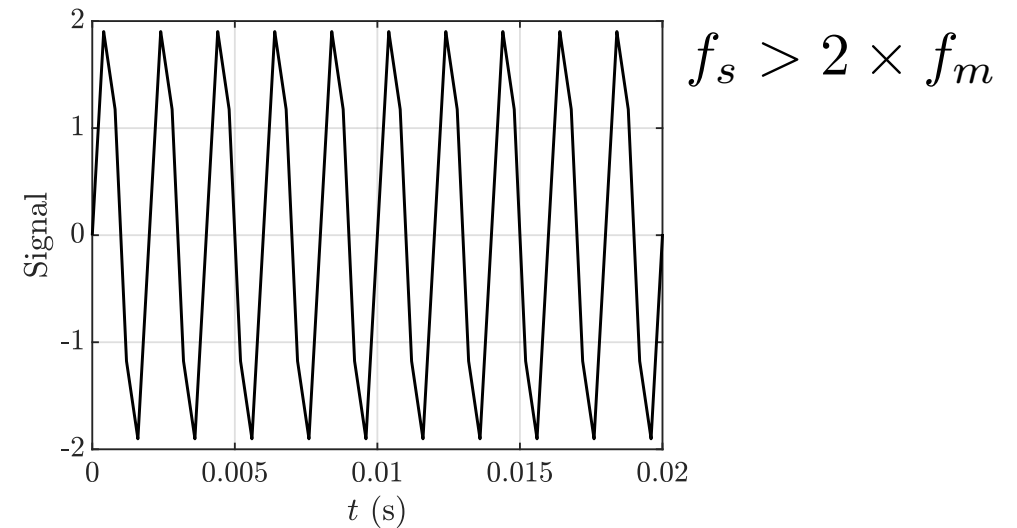
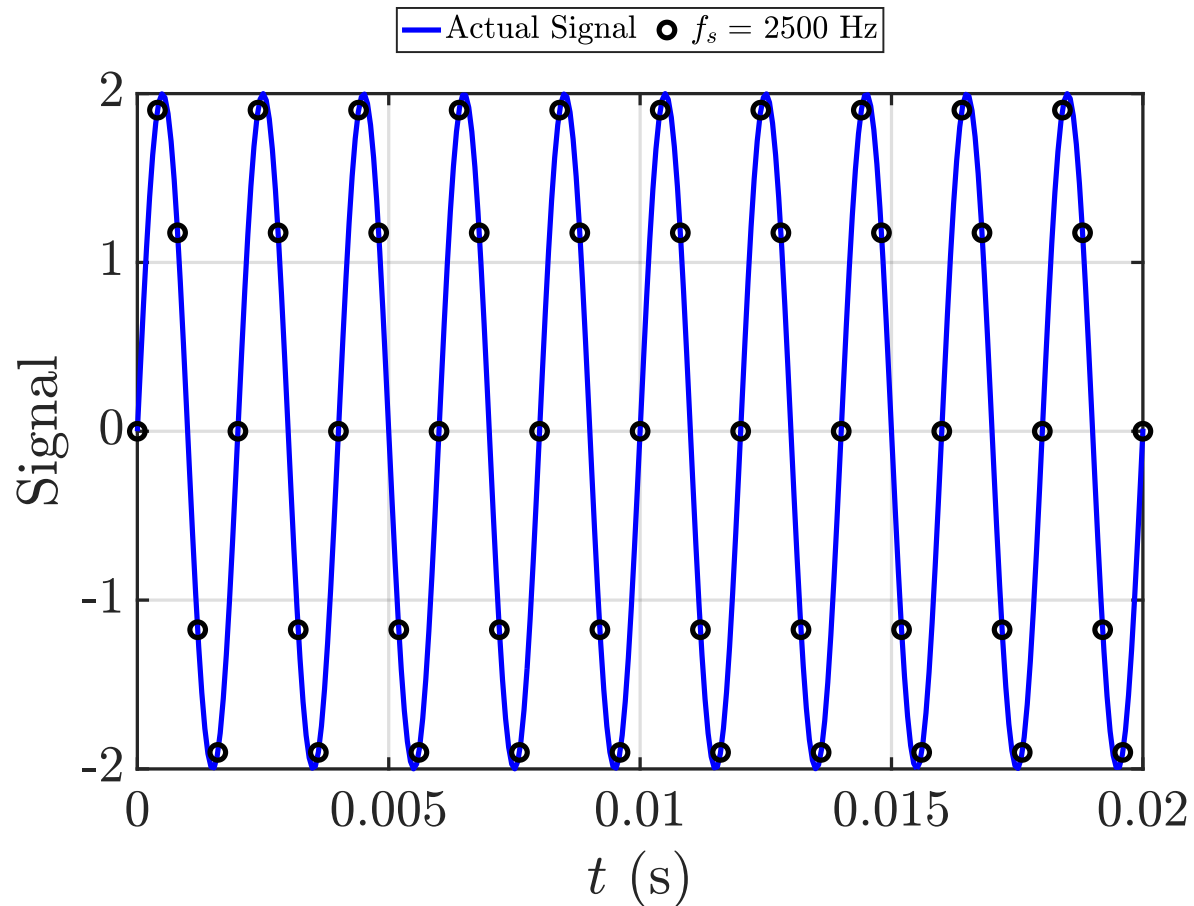
$$f_m = 500 \text{ Hz}$$



# Sampling theorem

Nyquist-Shannon criterion:  $f_s > 2 \times f_m$

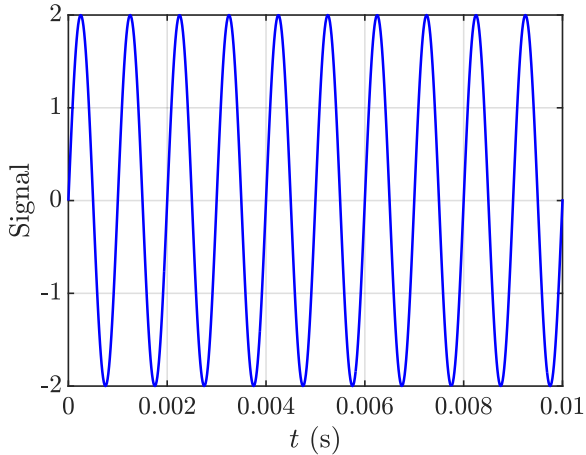
$$f_m = 500 \text{ Hz}$$



We got the frequency correct, but the amplitude?

In practice, because of noise etc., a sampling frequency **larger than ten times** is better

# Spectral analysis



## Fourier transform

Time domain  $\rightarrow$  Frequency domain

$$X_c(f) = \int_{-\infty}^{+\infty} x_c(t) e^{-i2\pi ft} dt$$

Frequency domain Time domain

For discrete signals

## Discrete Fourier transform

$$X(\omega) = \sum_{n=0}^{N-1} x(t) e^{-i\omega t}$$

$$t = n \times \frac{1}{f_s}$$

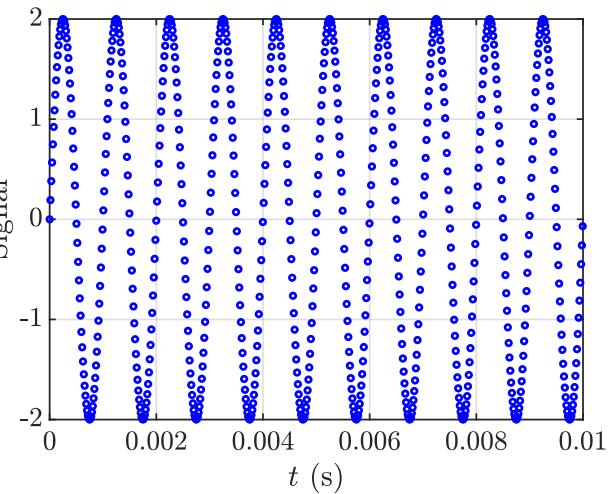
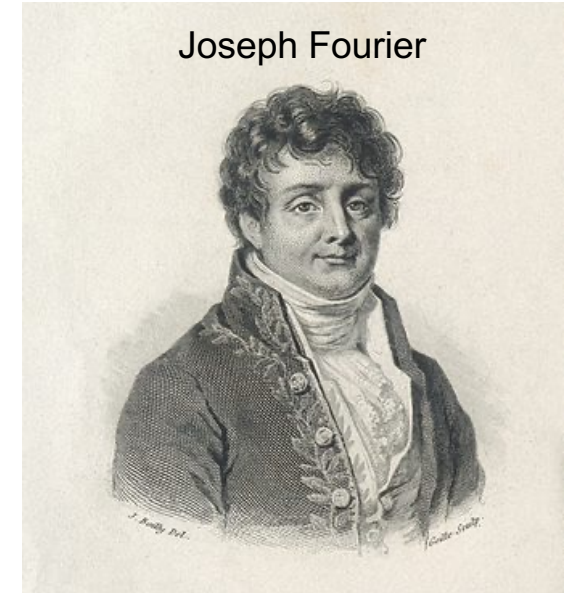
N – No. of points

## Fast Fourier transform algorithm

Cooley – Tukey algorithm

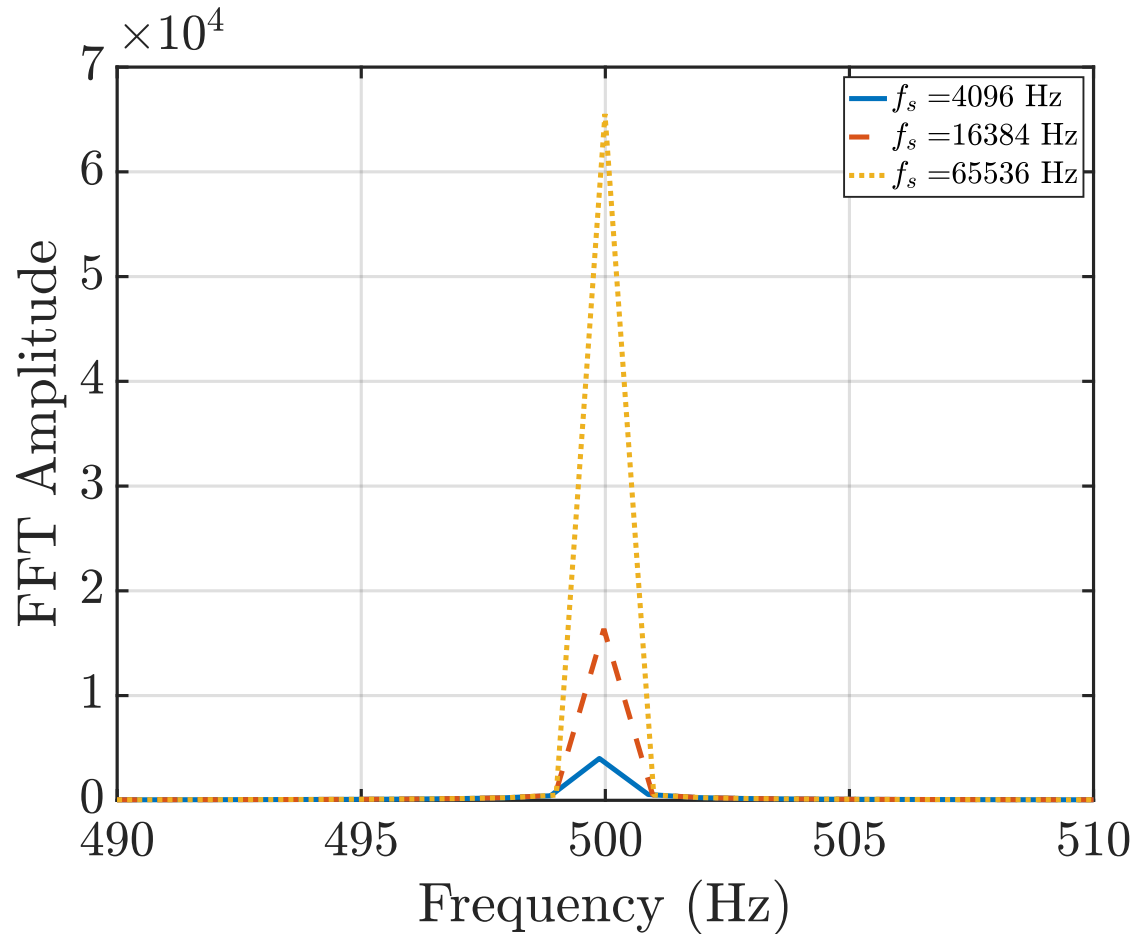
For numerical efficiency, no. of samples is a power of 2

Joseph Fourier



# Spectral analysis – FFT dependance on sampling rate

Signal :  $x(t) = 2 \sin(2\pi \cdot 500 \cdot t)$



$$X(\omega) = \sum_{n=0}^{N-1} x(t)e^{-i\omega t}$$

Amplitude depends on N  
number of points in DFT

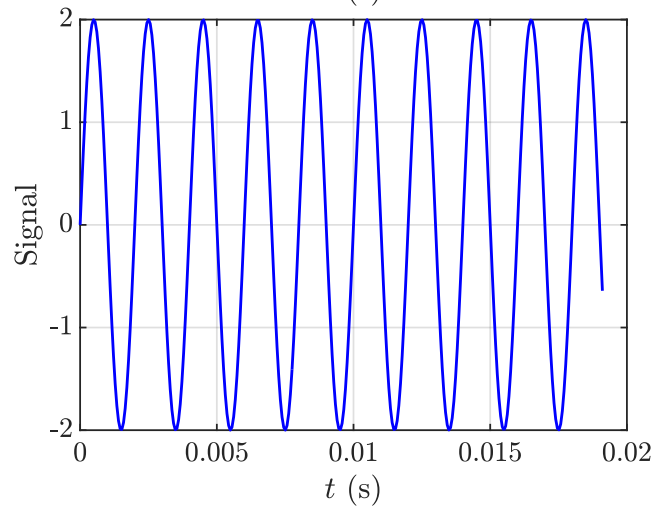
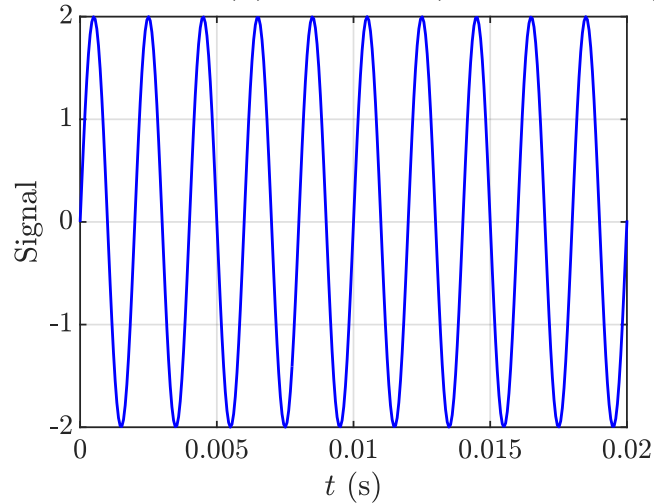


Normalize by frequency  
resolution

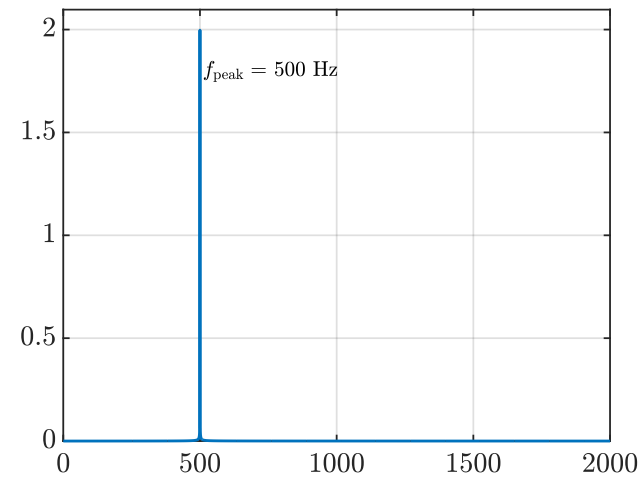


# Spectral analysis – spectral leakage

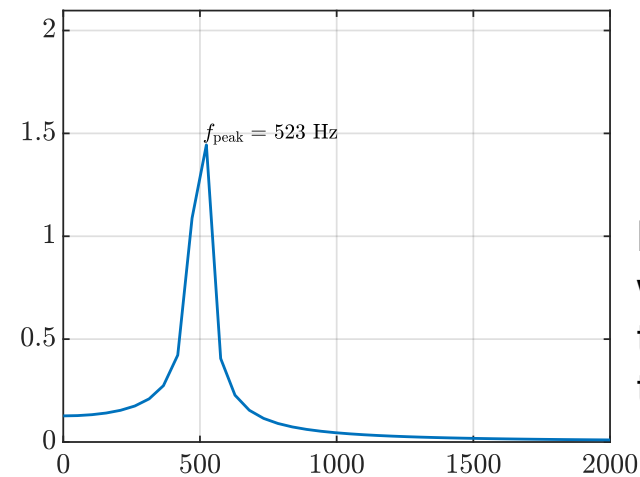
Signal :  $x(t) = 2 \sin(2\pi \cdot 500 \cdot t)$



**Fast Fourier transform**



DFT assumes the signal is periodic



Multiplying the signal by windows of finite length in which the amplitude gradually reduces to zero at the edges

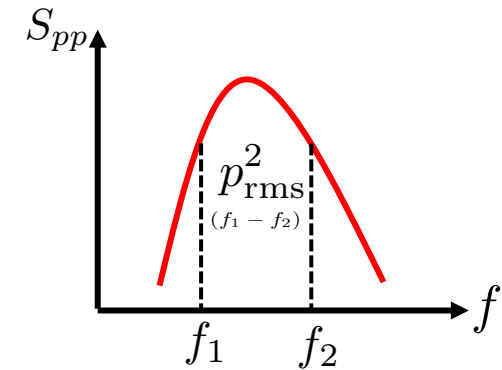
This can be avoided by windowing

# Power spectral density

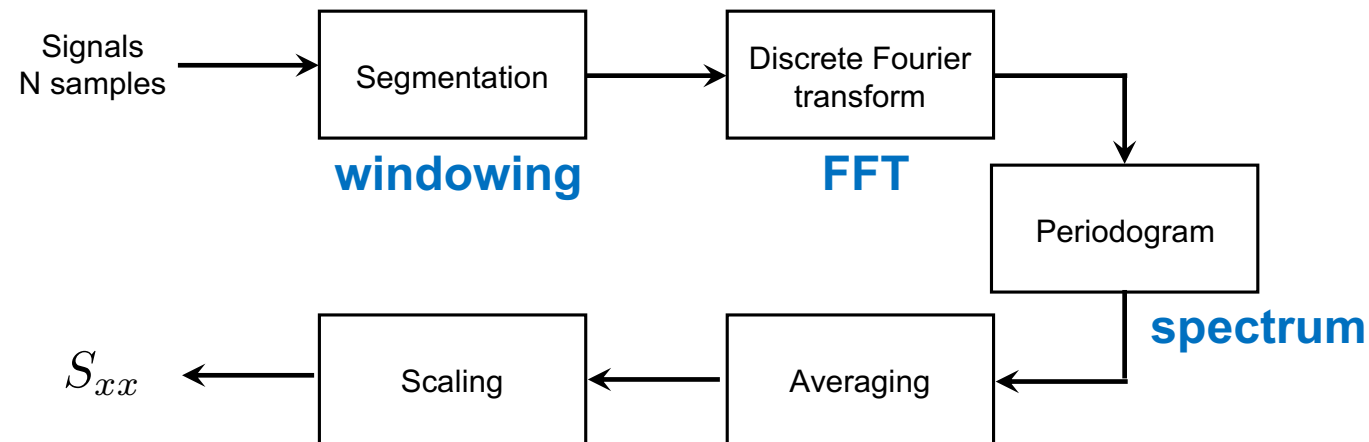
## Power spectral density

mean square value per unit frequency

$$p_{\text{rms}}^2 = \int S_{pp}(f) df$$

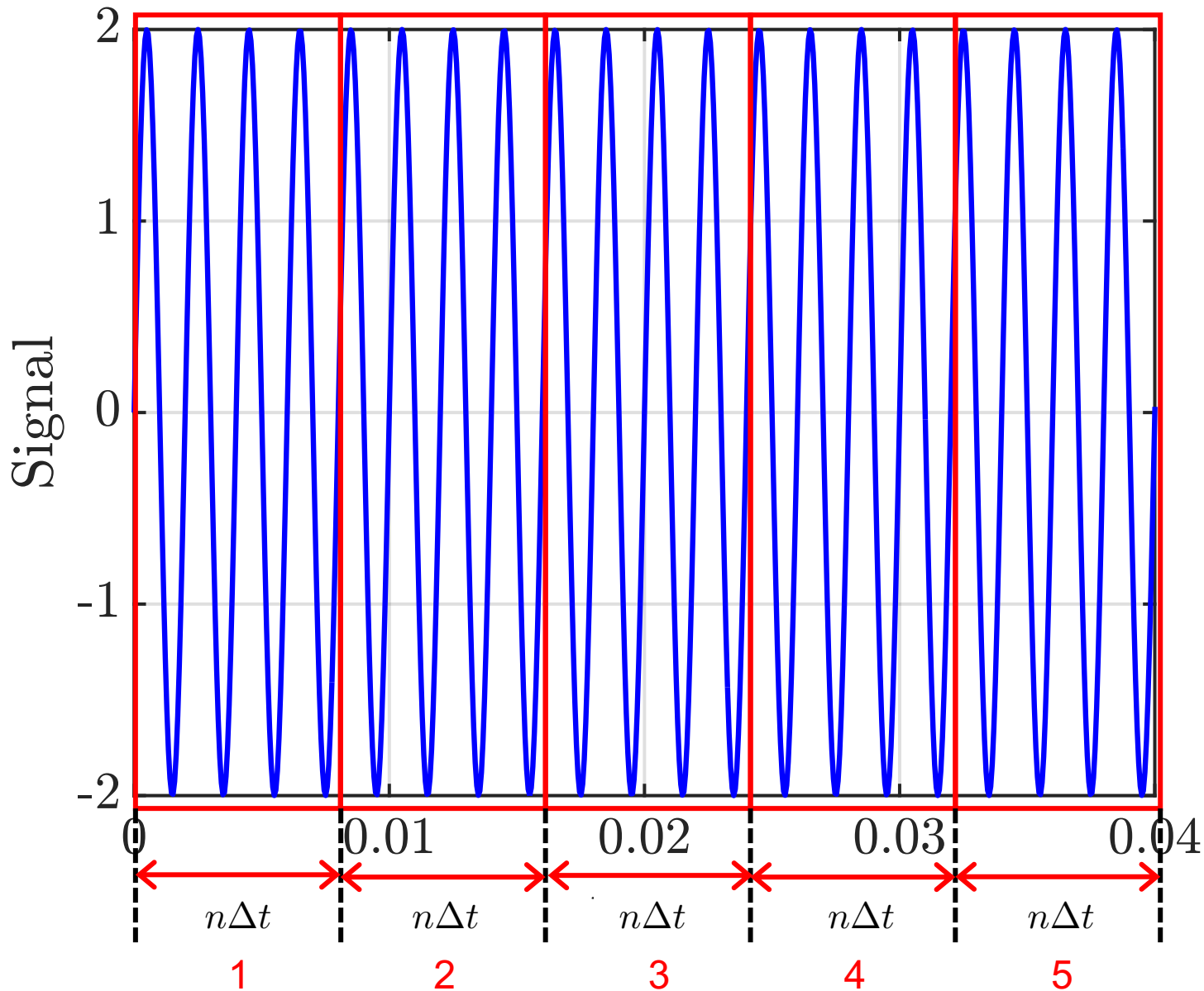


## Estimating power spectral density using Welch's periodogram technique

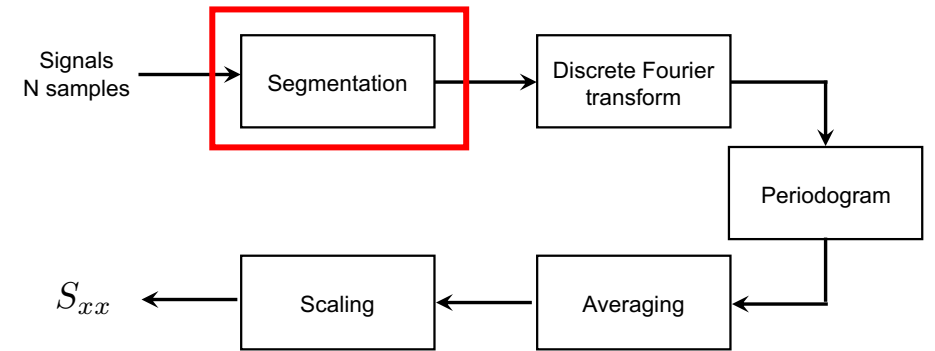


Reproduced from Denis Veynante, Survey of signal processing techniques

# Spectral analysis – Windowing



Signal :  $x(t) = 2 \sin(2\pi \cdot 500 \cdot t)$



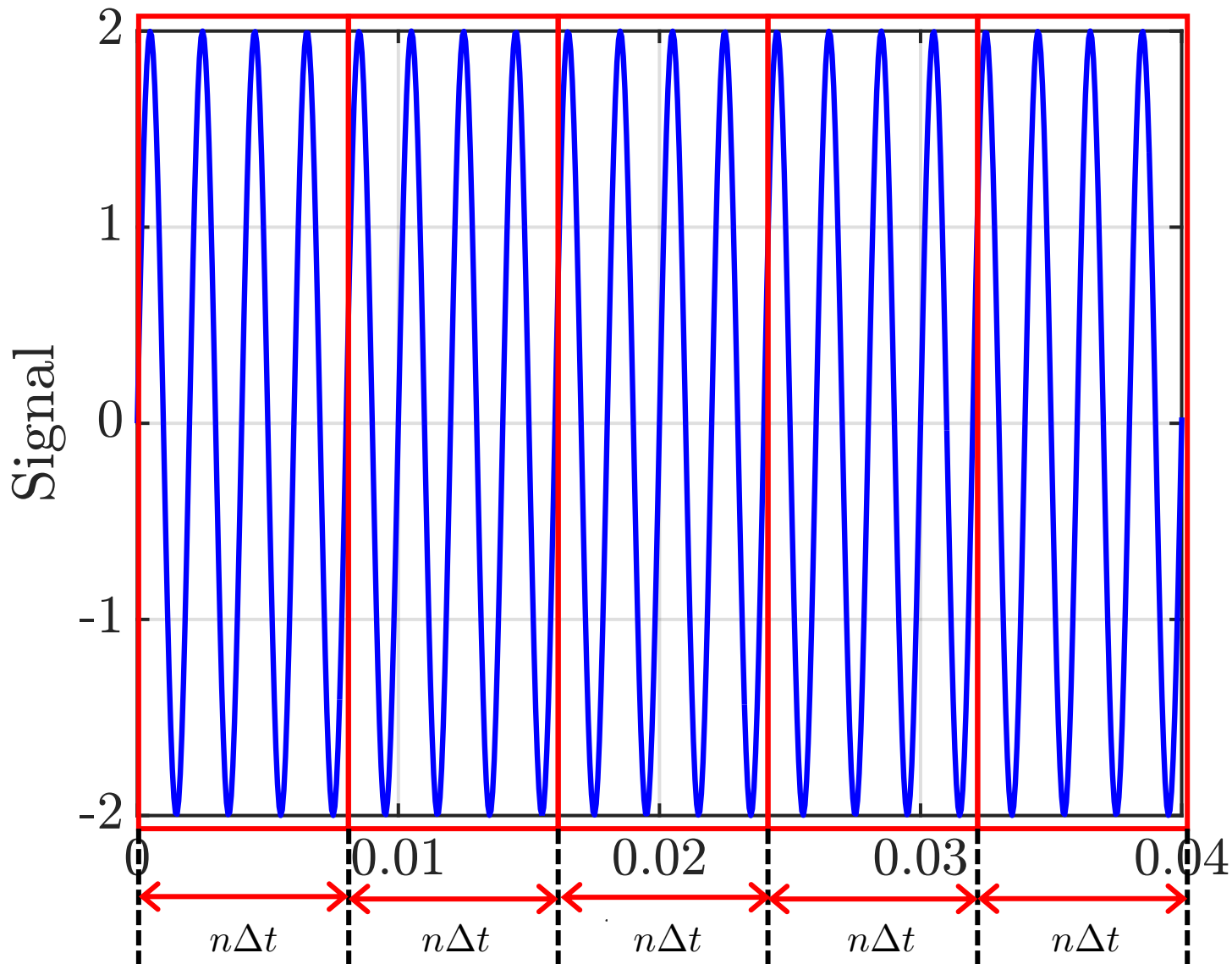
No. of windows:  $M$

No. of points/window:  $n$

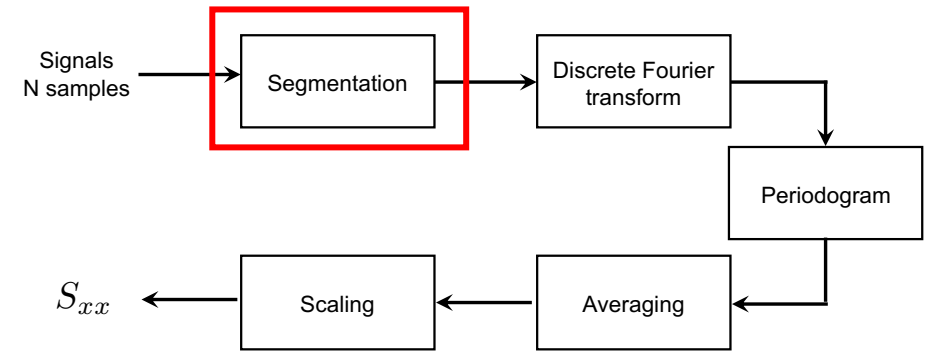
$$N = M \times n$$

Different types of windowing functions exist  
Hanning, Hamming etc.  
Choice depends on the application

# Spectral density using Welch's periodogram approach



Signal :  $x(t) = 2 \sin(2\pi \cdot 500 \cdot t)$



No. of windows:  $M$

No. samples/window:  $n$

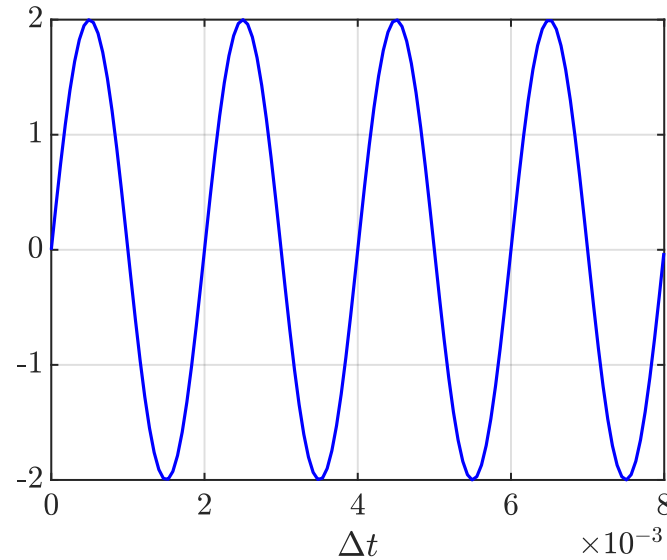
$$N = M \times n$$

Frequency resolution  $\Delta f = \frac{1}{n\Delta t}$

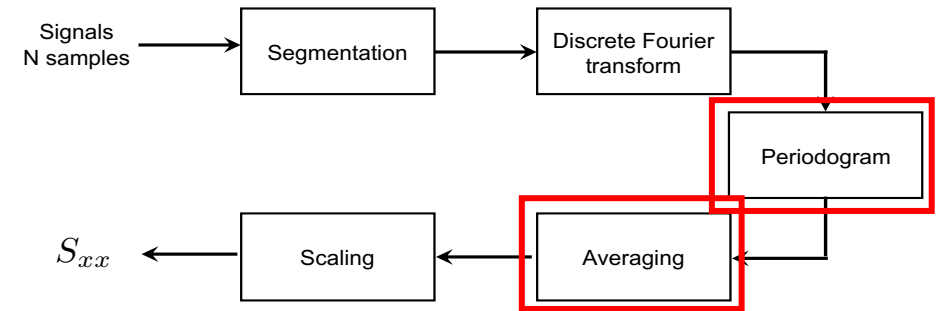
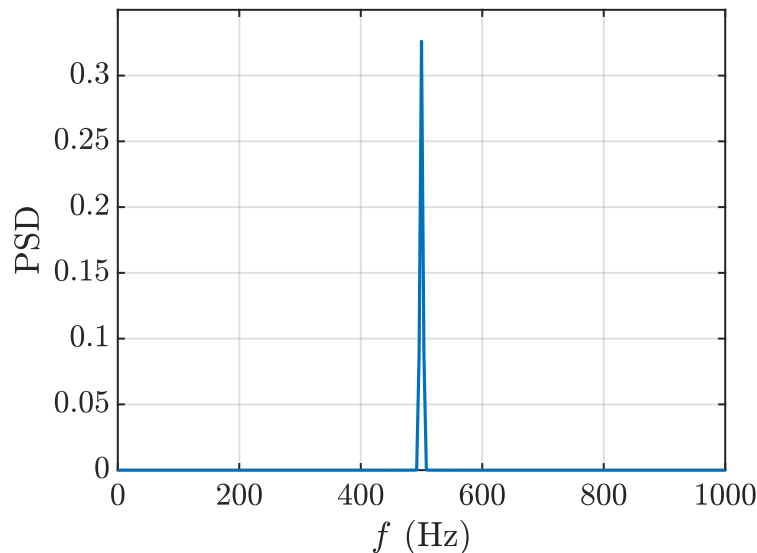
$\Delta t$  is the time step

# Spectral density using Welch's periodogram approach

Windowed signal



Periodogram



In the next step, averaging of periodograms from different windows is averaged

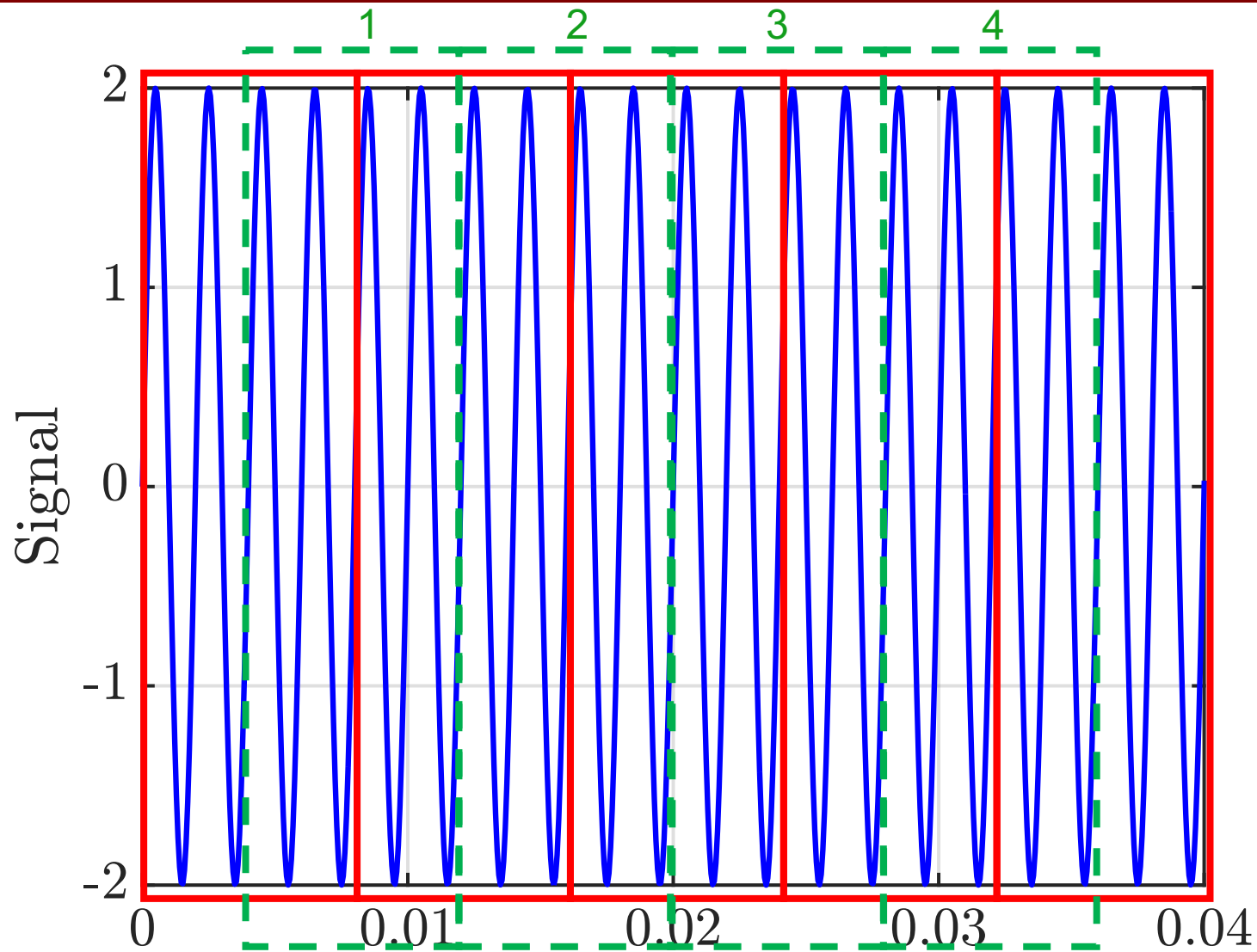
Practical signals have noise – averaging operation smoothens the effect of noise

Increasing  $M$  improves statistical stability and provides better spectral estimates

Frequency resolution decreases

$$\Delta f = \frac{1}{n\Delta t}$$

# Spectral estimate – overlapping



Improving statistical stability without compromising frequency resolution

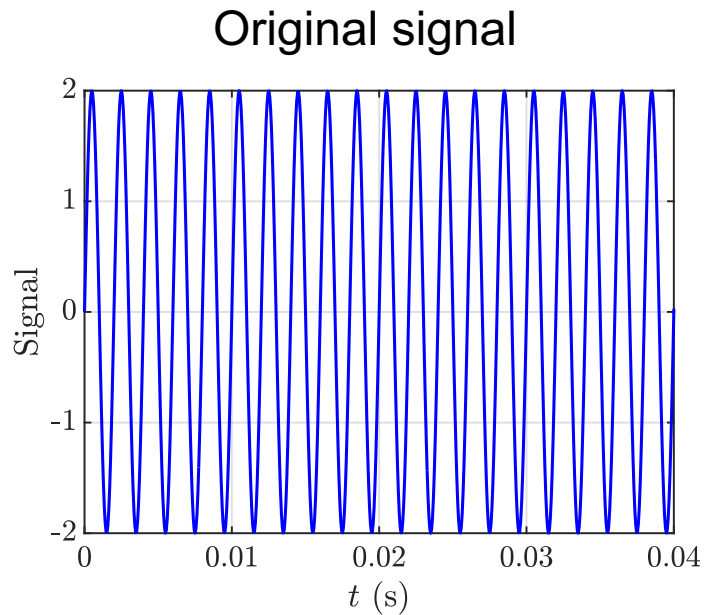
## Overlapping

No. of original **windows** = 5

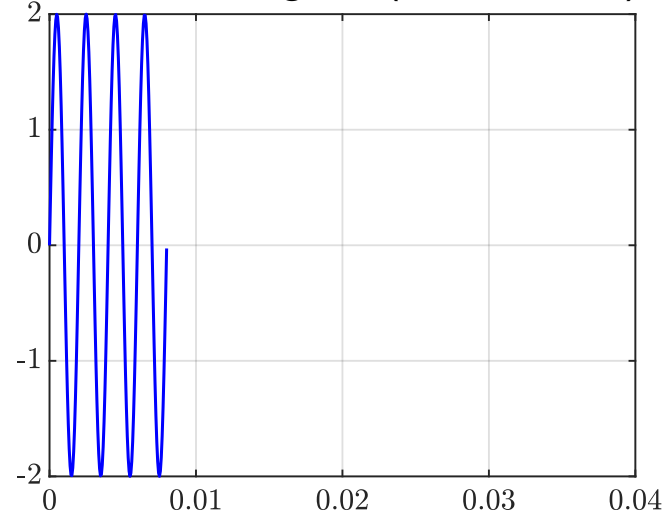
with 50% **overlap**

Total no. of windows = 9

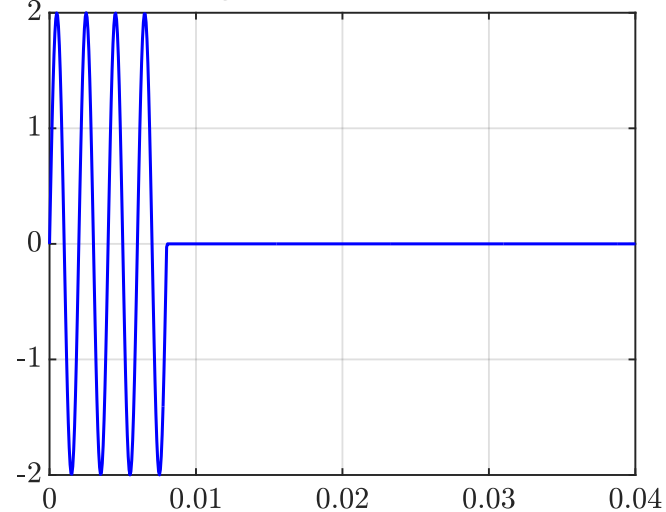
# Spectral estimate – Zero-padding



windowed signal (5 windows)



windowed signal after zero-padding



Artificially increasing  
frequency resolution

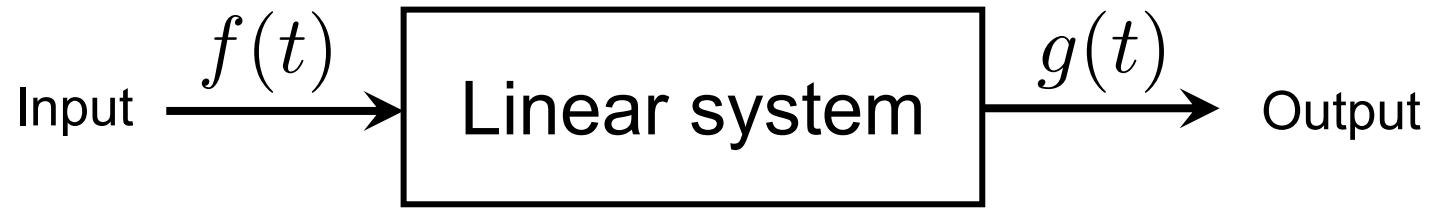
$$\Delta f = \frac{1}{n\Delta t}$$

$n$  is increased after zero-padding

No actual information in zero-  
padded data

It is not exactly resolution that  
increases but it is rather the  
frequency precision

# Transfer function



$$S_{fg}(\omega) = H(j\omega)S_{ff}(\omega)$$

Cross-power spectral density ←      ↓      → Power spectral density  
Transfer function

**Gives a gain and a phase**



# Coherence function

Useful when the output is noisy, which is the case for all practical signals

Degree of dissimilarity of two signals

**Coherence function**

$$\gamma_{fg}^2(\omega) = \frac{|S_{fg}(\omega)|^2}{S_{ff}(\omega)S_{gg}(\omega)}$$

f(t) : input  
g(t) : output

$$0 \leq \gamma_{fg}^2 \leq 1$$

If there is no source of noise corrupting output, coherence is 1

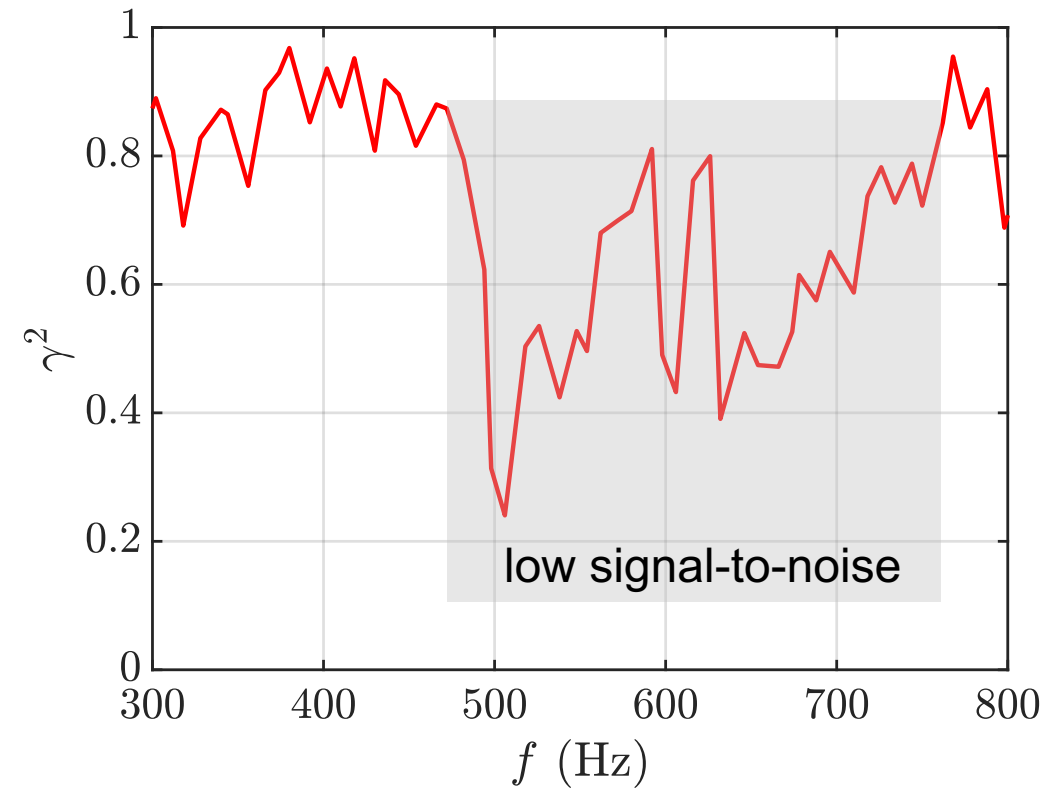
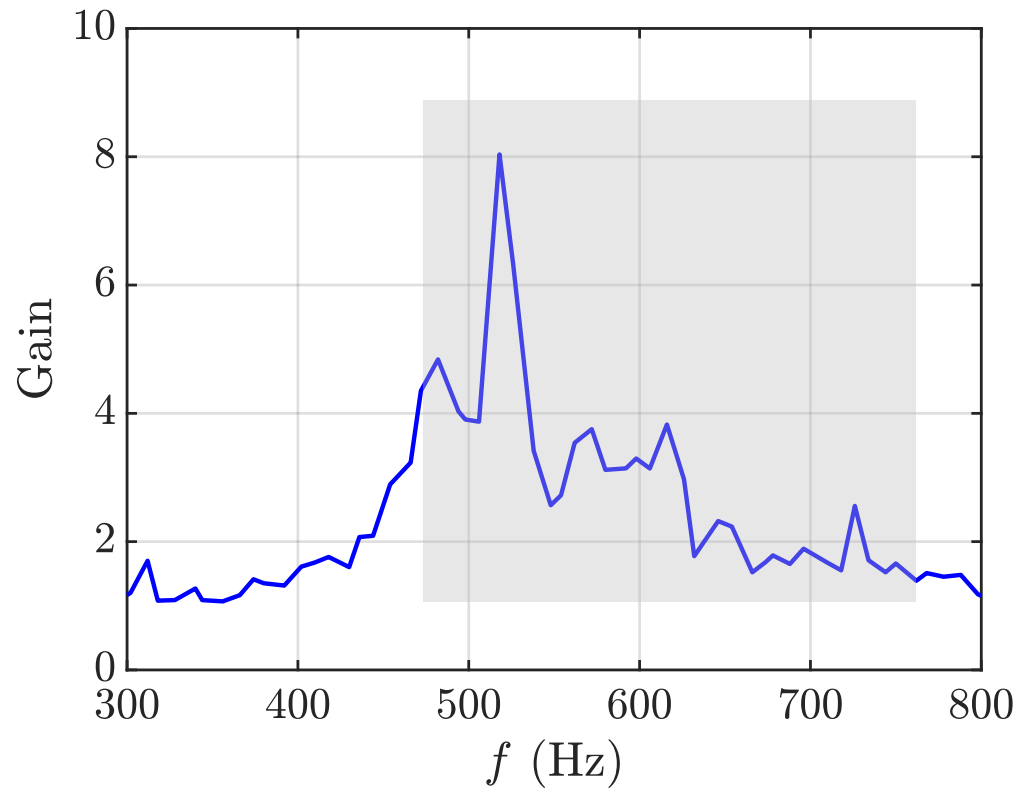
**Signal-to-noise ratio**

$$\frac{S_{gg}(\omega)}{S_{nn}(\omega)} = \frac{\gamma_{fg}^2(\omega)}{1 - \gamma_{fg}^2(\omega)}$$

# Coherence function

**Coherence function**

$$\gamma_{fg}^2(\omega) = \frac{|S_{fg}(\omega)|^2}{S_{ff}(\omega)S_{gg}(\omega)}$$



# Flame transfer function



$$\mathcal{F}(\omega) = \frac{\dot{Q}'(\omega)/\bar{\dot{Q}}}{u'(\omega)/\bar{u}}$$

In premixed or quasi-premixed mode of combustion

Linear description cannot explain nonlinear features limit cycle oscillations, mode switching etc. that are commonly found in a system exhibiting a thermoacoustic instability

# Flame describing function

**Flame transfer function**

$$\mathcal{F}(\omega) = \frac{\dot{Q}'(\omega)/\overline{\dot{Q}}}{u'(\omega)/\bar{u}}$$

**Flame describing function**<sup>1,2</sup> is an extension of flame transfer function, and it represents the flame response depending on the frequency and amplitude of the incoming velocity perturbations

**Flame describing function**

$$\mathcal{F}(\omega, |u'|) = \frac{\dot{Q}'(\omega, |u'|)/\overline{\dot{Q}}}{u'(\omega)/\bar{u}}$$

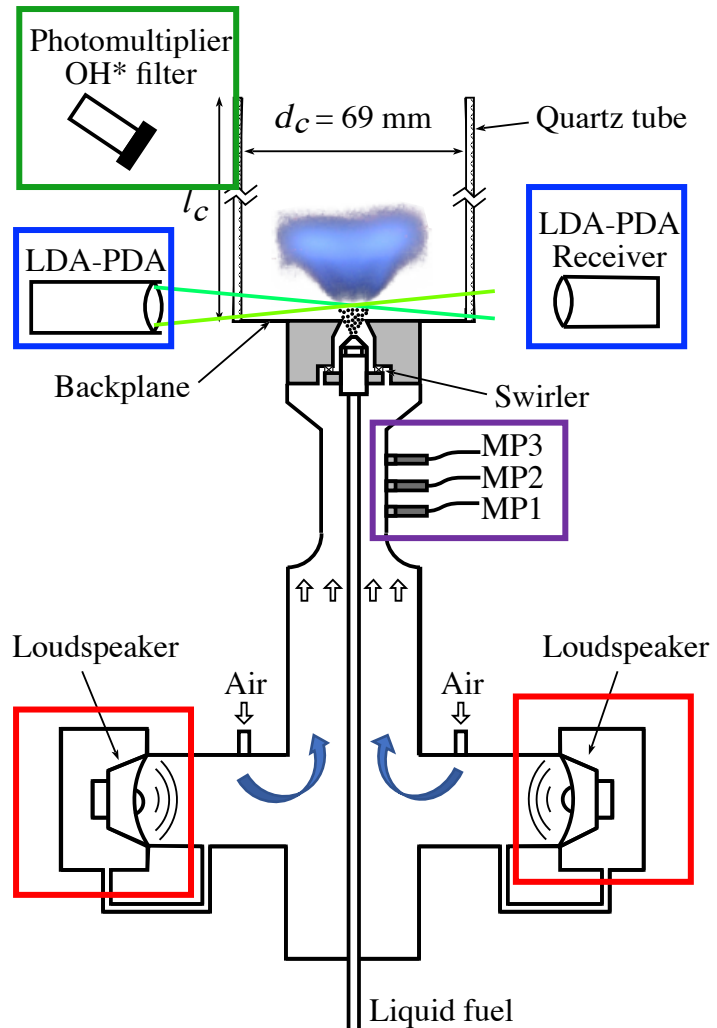
Gain & phase information can be obtained

[1] **A. P. Dowling** (1997) *J. Fluid Mechanics*, 346.

Non-linear self-excited oscillations of a ducted flame

[2] **N. Noiray, D. Durox, T. Schuller, and S. Candel** (2008) *J. Fluid Mechanics*, 615. A unified framework for nonlinear combustion instability analysis based on the flame describing function

# SICCA-Spray combustor – experiments



## Measuring equipments

Light intensity fluctuations – Photomultiplier with OH\* filter  
OH\* chemiluminescence at 308 nm  
Approximated to heat release rate fluctuations

Velocity fluctuations (injector exit) – Laser Doppler anemometry  
Measures mean & fluctuating velocity

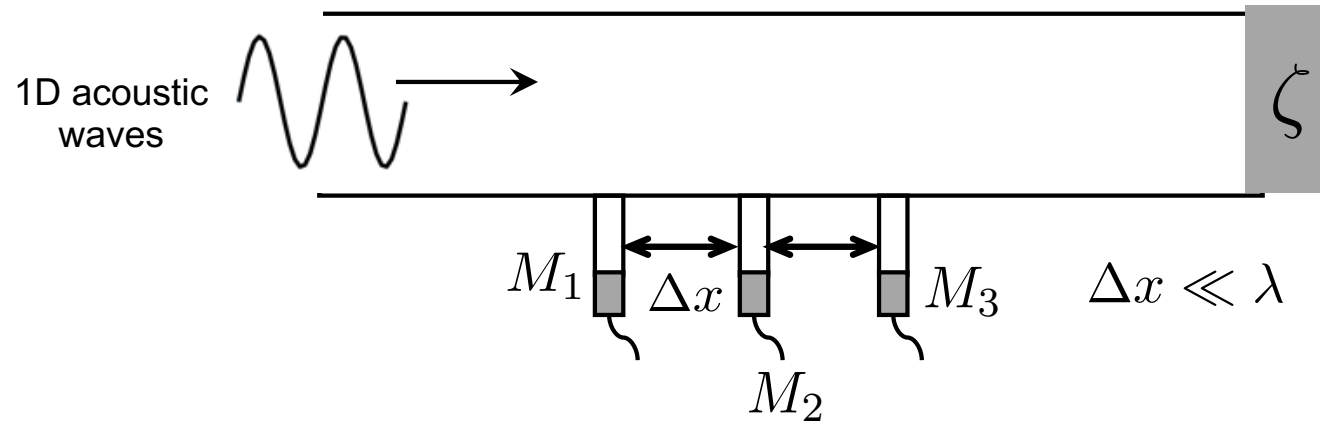
Velocity fluctuations (plenum) – Multi-microphone method  
Calculate acoustic velocity from acoustic pressure signals

## Input signal

Flow modulation – loudspeakers  
Pulsations produced at different amplitude and frequencies

# Multi-microphone method

Recall from POLKA's 2<sup>nd</sup> scientific workshop – Prof. Hans Boden



Two microphones are mostly sufficient  
But with third microphone, fidelity can  
be improved

**Acoustic momentum  
equation**

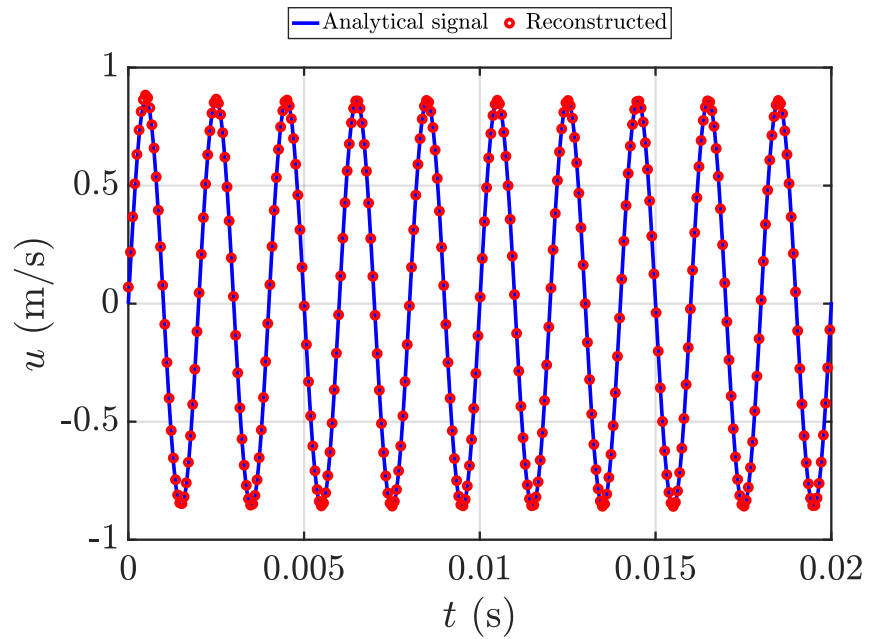
$$\frac{\partial u}{\partial t} = \frac{-1}{\rho_0} \frac{\partial p}{\partial x}$$

$$u \simeq \frac{-i}{\rho_0 \omega \Delta x} (p(x_2) - p(x_1))$$

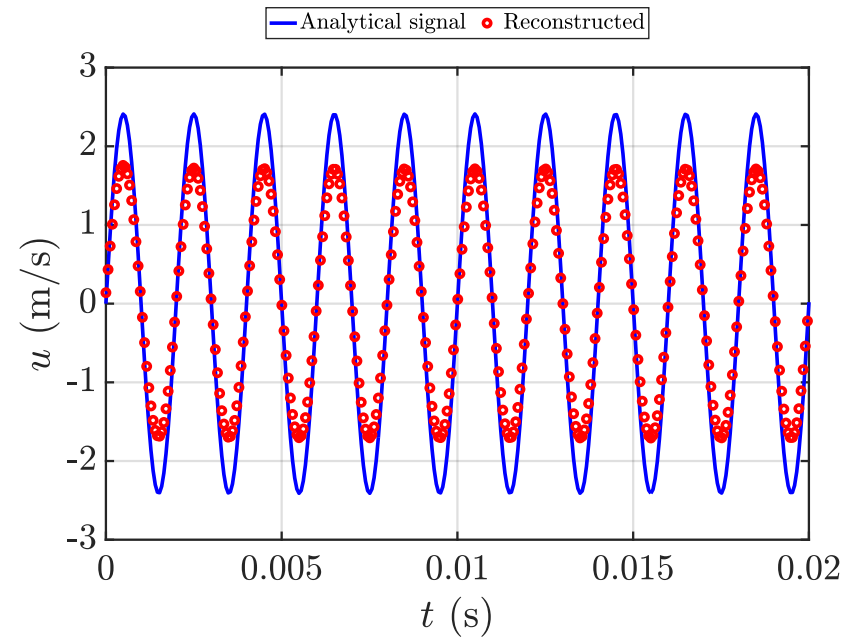
Analytical pressure signals using Hilbert  
transform is necessary

# Multi-microphone method

$$\Delta x \ll \lambda$$



$$\Delta x \simeq \lambda$$



# SICCA-Spray combustor – experiments

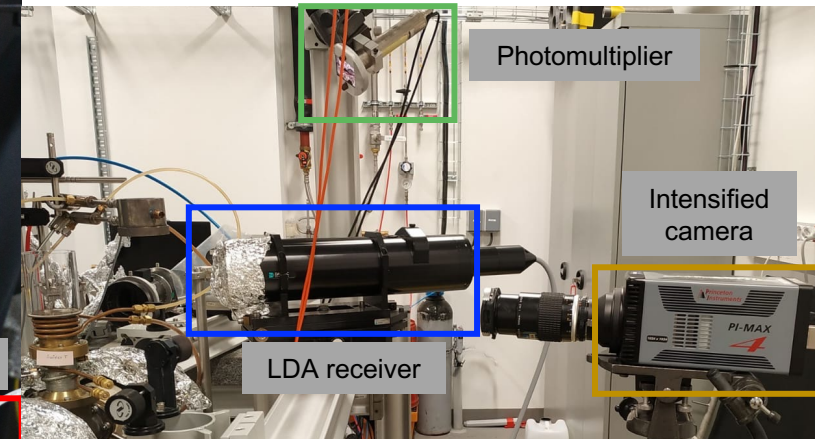
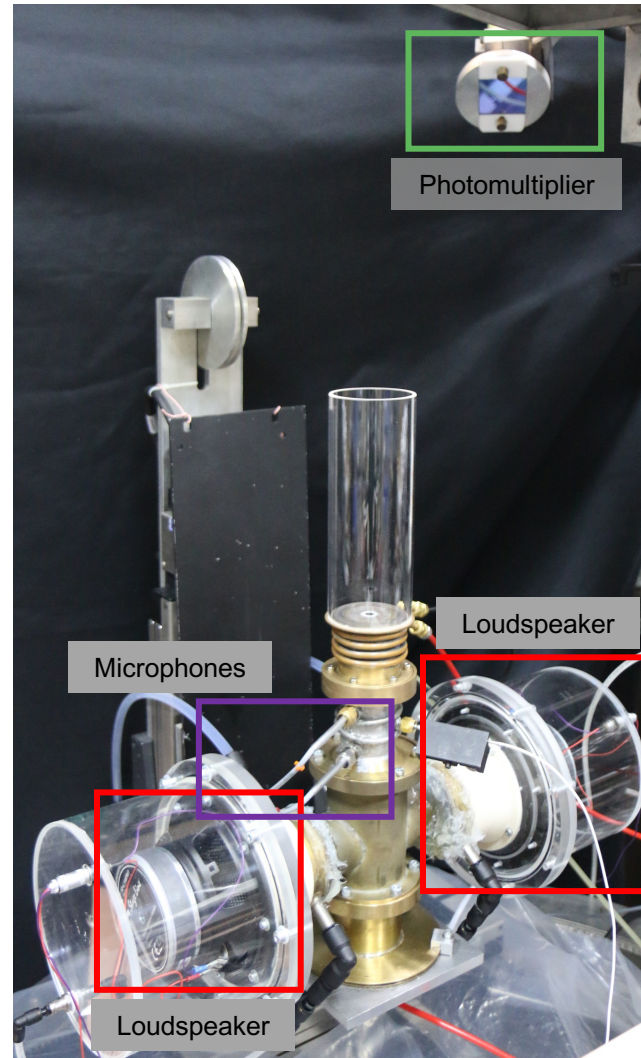
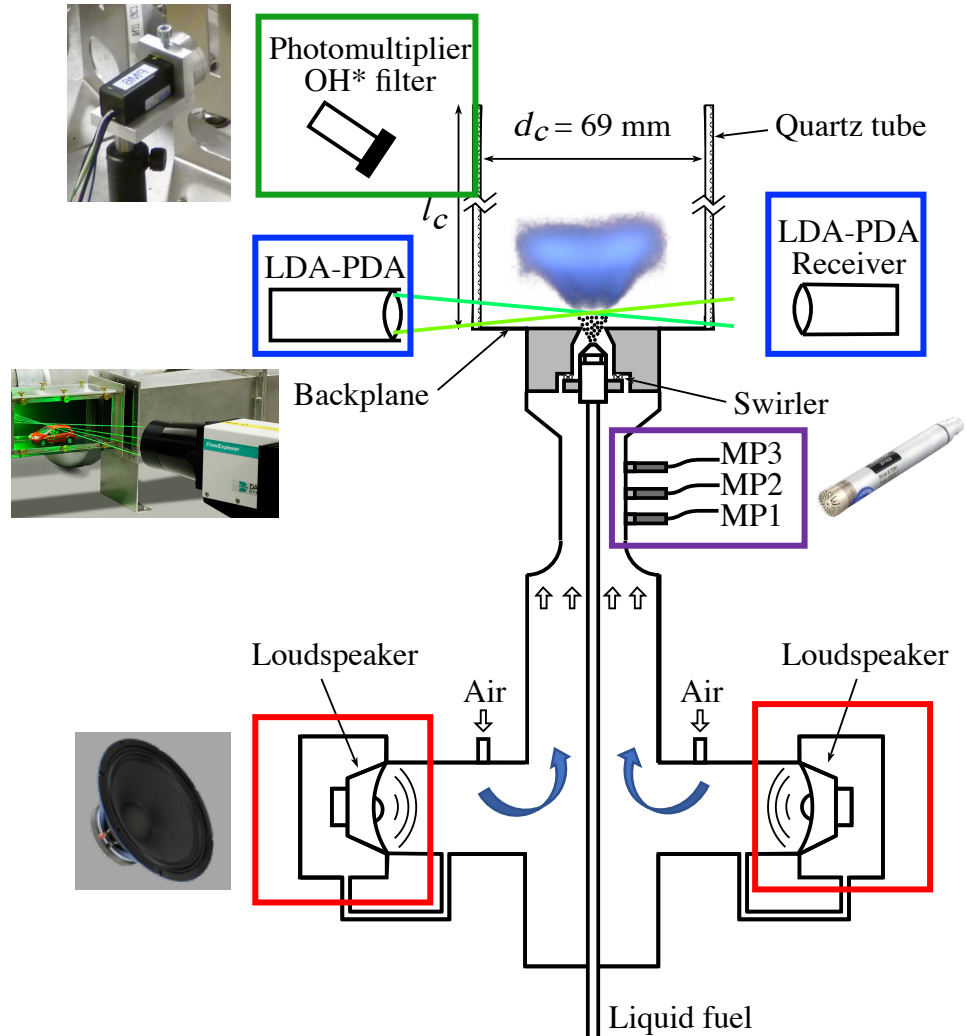
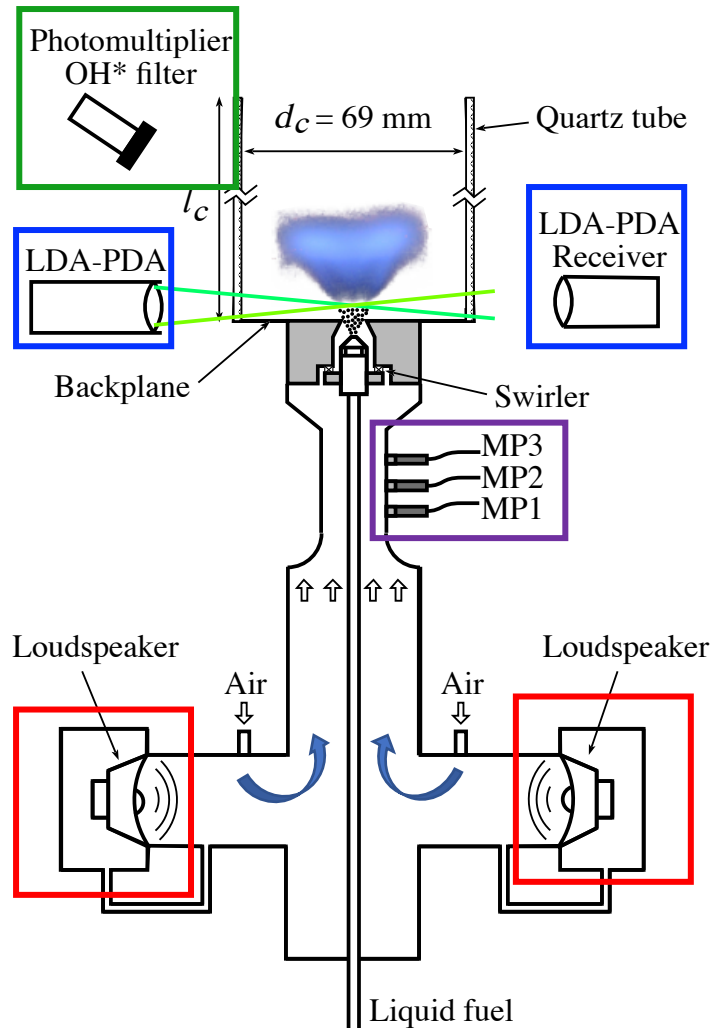


Image sources  
 Microphone : Bruel & Kjaer  
 Loudspeaker: Monacor  
 LDA system: Dantec



# SICCA-Spray combustor – experiments



In the input .mat file, look for

Photomultiplier signal: **PM** (V)

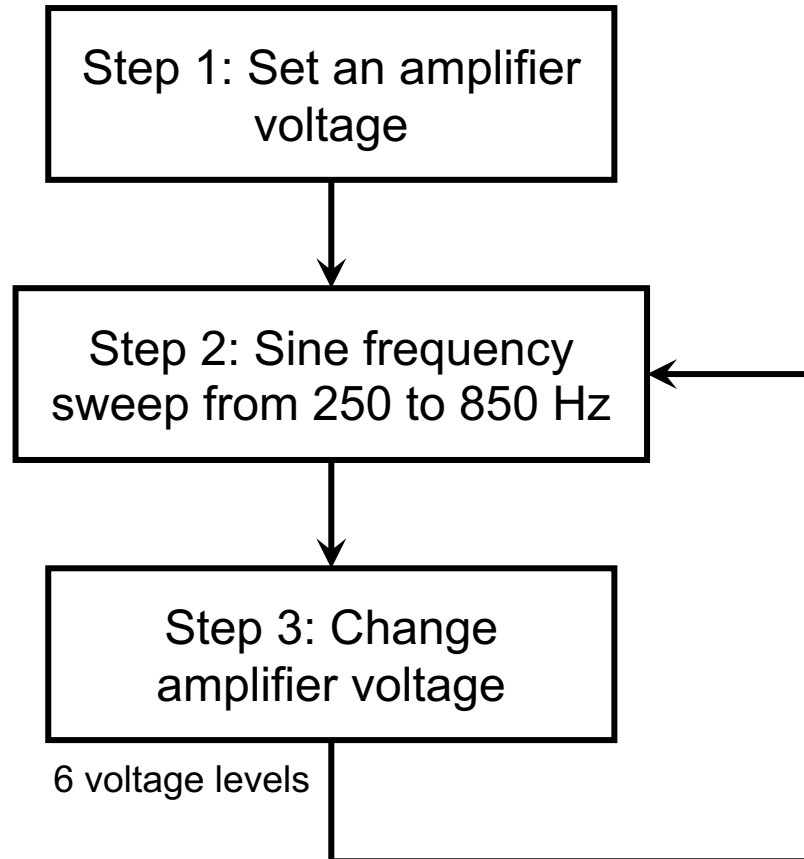
LDA signal (chamber velocity): **u** (m/s)

Microphones (plenum) – only two  
MP1: **Pr2** - bottom mic (Pa)  
MP3: **Pr4** - top mic (Pa)

Input signal to loudspeakers  
Signal generator: **Gene** (V)  
Generator frequency: **refFreq** (Hz)  
Reference signal – sinusoidal

# SICCA-Spray combustor – acquisition

## How is the measurement carried out?



Amplifier voltages:

500, 1000, 1500, 2000, 2500, 3000 mV

Frequency ramp range: 250 Hz – 850 Hz

Frequency ramp rate: 4.5 Hz/s

Measurement acquired in blocks of 2s

Measurement obtained at different frequencies and amplitudes

# Matlab exercise

- Open the script: Exercise\**POLKA\_exercise.m**
- Data file: Exercise\Data
- Section that need to be completed are commented as “**TO BE COMPLETED**”
- For determining acoustic velocity **VelAcoustic.m** function should be completed

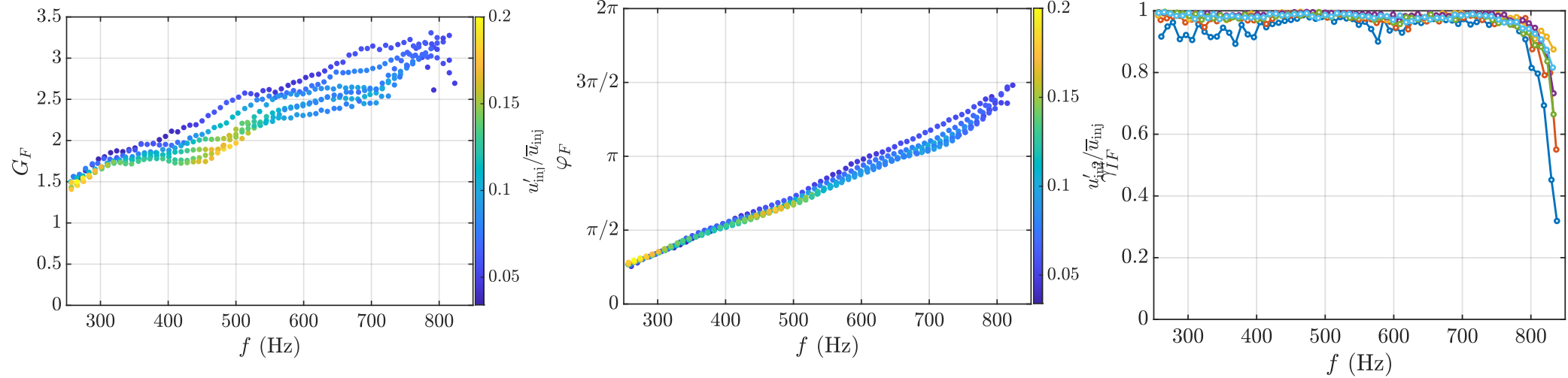
# Calculating flame describing function

You are given:

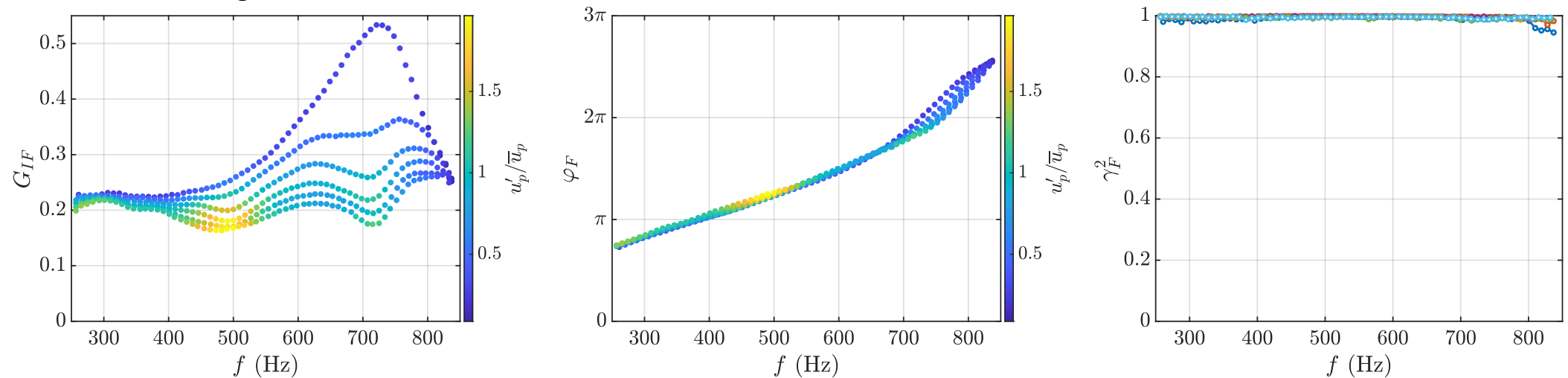
- Microphone signals: **P2**, **P4** (Pa) in plenum
- Velocity signal: **u** (m/s) at injector exit
- Photomultiplier signal: **PM** (V)
- Reference signal from generator: **Gene** (V)
- Frequency of reference signal: **refFreq** (Hz)

# Results

## Flame describing function



## Injector + flame describing function



# Step 1: Bandpass filtering

- Remove noise from measured signals
  - **Command:** bandpass
  - Specify **lower cut-off and upper cut-off** frequencies (already given)
  - Specify the **sampling rate** (Fs)
  - **Variables:** uBlock, PMBlock, Pr2Block, Pr4Block

```
% ----- Bandpass filtering (TO BE COMPLETED) ----- %
% use the reference frequency for bandpass filtering microphone, velocity
% and PM signals: filter 5% around the reference frequency
% Matlab command: bandpass

freqFilterLower = refFreq(i) - 5/100*refFreq(i);
freqFilterHigher = refFreq(i) + 5/100*refFreq(i);

% calculating mean values for normalizing transfer functions
uMeanInjector = mean(uBlock);
uBulkPlenum = 2.7; % m/s predetermined
PMMean = mean(PMBlock);

% use variables: uBlock for velocity
%                 PMBlock for light intensity
%                 Pr2Block for lower plenum microphone
%                 Pr4Block for upper plenum microphone

uFilt = bandpass(uBlock,[freqFilterLower,freqFilterHigher],Fs); % filterin
PMFilt = bandpass(PMBlock,[freqFilterLower,freqFilterHigher],Fs); % filter
Pr2Filt = bandpass(Pr2Block,[freqFilterLower,freqFilterHigher],Fs); % filt
Pr4Filt = bandpass(Pr4Block,[freqFilterLower,freqFilterHigher],Fs); % filt
```

# Step 2: Finding acoustic velocity in the plenum

Fill in **VelAcoustic.m** function

$$u \simeq \frac{-i}{\rho_0 \omega \Delta x} \left( \underbrace{p(x_2)}_{\text{top mic}} - \underbrace{p(x_1)}_{\text{bottom mic}} \right)$$

$\downarrow$   $\quad \quad \quad \underbrace{\hspace{10em}}$   
dX  $\quad \quad \quad$  dP Use **hilbertRev** function

$$u = \frac{\text{imag}(\Delta P)}{2\pi f \rho \Delta x}$$

```
function[u] = VelAcoustic(pressure_bottom,pressure_top,dX,rho,freq)

% 1. Find pressure difference
% and determine the Hilbert transform using hilbertRev function

dP = hilbertRev(pressure_top-pressure_bottom); % hilbert transform w

% 2. Determine u using the formula
% dX - distance between microphones
% rho - density
% freq - frequency of the microphone signals
u = imag(dP) / (rho*dX*2*pi*freq);

end
```

# Step 3: Calculating transfer functions

1. Calculate cross-power spectral density between velocity & heat release rate signals

- **Command:** cpsd
- **Variables:** input, output
- **Welch's parameters:** dataWindow, noverlap, n\_fft, Fs

2. Calculate power spectral density of velocity signals

- **Command:** pwelch
- **Variables:** input
- **Welch's parameters:** dataWindow, noverlap, n\_fft, Fs

```
% Transfer function - injector + flame

input = uPlenum./uBulkPlenum; % this is relative velocity fluctuations in
output = PMFilt./PMMean; % this is relative light intensity fluctuations f

% CPSD between input and output
[cpsdAmpIF,cpsdFreqIF] = cpsd(input,output,dataWindow,noverlap,n_fft,Fs);
[~,indexForMax] = max(abs(cpsdAmpIF));
cpsdPeakIF = cpsdAmpIF(indexForMax);
freqIF(i,vloop) = cpsdFreqIF(indexForMax);

% PSD of input
[psdAmpIF,psdFreqIF] = pwelch(input,dataWindow,noverlap,n_fft,Fs);
[~,indexForMax] = max(psdAmpIF);
psdPeakIF = psdAmpIF(indexForMax);
```

First do it for plenum signals → to obtain injector + flame transfer function

Then do it for chamber signals → to obtain flame transfer function



# Step 3: Calculating transfer functions

- To calculation gain
  - **Command:** abs(cpsd/psd)
  - **Variables:** cpsdPeakIF, psdPeakIF for injector + flame transfer function
  - **Variables:** cpsdPeakF, psdPeakF for flame transfer function

- To calculation phase

```
gainIF(i,vloop) = abs(cpsdPeakIF/psdPeakIF);  
phaseIF(i,vloop) = angle(cpsdPeakIF);
```

- **Command:** angle(cpsd)
- **Variables:** cpsdPeakIF for injector + flame transfer function
- **Variables:** cpsdPeakF for flame transfer function

- To calculation coherence function

- **Command:** abs(cpsd)^2/(psd-input \* psd-output)
- **Variables:** cpsdPeakIF, psdPeakIF, psdPeakOut
- **Variables:** cpsdPeakF, psdPeakF, psdPeakOut

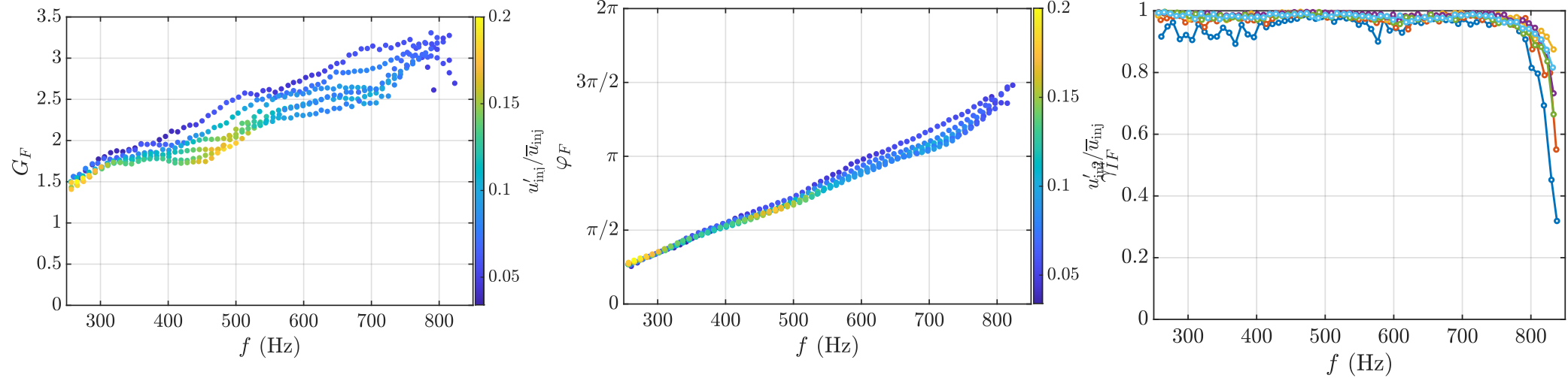
$$\gamma_{fg}^2(\omega) = \frac{|S_{fg}(\omega)|^2}{S_{ff}(\omega)S_{gg}(\omega)}$$

```
coherenceIF(i,vloop) = abs(cpsdPeakIF)^2/(psdPeakIF*psdPeakOut);
```

f(t) : input; g(t) : output

# Results

## Flame describing function



## Injector + flame describing function

